

$$\begin{aligned}
&= \frac{\partial}{\partial x_j} \left[\overline{u'_i \left(\frac{\partial u'_i}{\partial x_j} + \frac{\partial u'_j}{\partial x_i} \right)} \right] - \overline{\frac{\partial u'_i}{\partial x_j} \frac{\partial u'_j}{\partial x_i}} + \underbrace{u'_i \frac{\partial}{\partial x_i} \left(\frac{\partial u'_j}{\partial x_j} \right)}_0 \\
\frac{u'_i u'_j \frac{\partial \bar{u}_i}{\partial x_j}} &= \frac{u'_j u'_i \frac{\partial \bar{u}_j}{\partial x_i}} = \frac{1}{2} \overline{u'_i u'_j \left(\frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right)} \\
- \overline{\left(\frac{\partial u'_i}{\partial x_j} \right)^2} - \overline{\frac{\partial u'_i}{\partial x_j} \frac{\partial u'_j}{\partial x_i}} &= -\frac{1}{2} \overline{\left(\frac{\partial u'_i}{\partial x_j} \right)^2} - \frac{1}{2} \overline{\left(\frac{\partial u'_j}{\partial x_i} \right)^2} - \overline{\frac{\partial u'_i}{\partial x_j} \frac{\partial u'_j}{\partial x_i}} = -\frac{1}{2} \overline{\left(\frac{\partial u'_i}{\partial x_j} + \frac{\partial u'_j}{\partial x_i} \right)^2}
\end{aligned}$$

を用いて

$$\begin{aligned}
&\frac{\partial}{\partial t} \left(\frac{1}{2} \overline{u'_i u'_i} \right) + \overline{u_j \frac{\partial}{\partial x_j} \left(\frac{1}{2} \overline{u'_i u'_i} \right)} \\
&= -\frac{\partial}{\partial x_j} \left(\frac{1}{\rho} \overline{u'_j p'} + \frac{1}{2} \overline{u'_i u'_i u'_j} - 2\nu \overline{u'_i s'_{ij}} \right) - \overline{u'_i u'_j s'_{ij}} - 2\nu \overline{s'_{ij} s'_{ij}} \quad (5.42)
\end{aligned}$$

となる。ただし、 $\overline{s'_{ij}}$, s'_{ij} は次のような定義である。

$$\overline{s'_{ij}} \equiv \frac{1}{2} \left(\frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right), \quad s'_{ij} \equiv \frac{1}{2} \left(\frac{\partial u'_i}{\partial x_j} + \frac{\partial u'_j}{\partial x_i} \right)$$

熱の分散 $\overline{\theta'^2}$ を支配する式

(5.34) に $2\theta'$ をかけて平均をとることによって求められる。

$$\frac{\partial \overline{\theta'^2}}{\partial t} + \overline{u_j \frac{\partial \overline{\theta'^2}}{\partial x_j}} = -\frac{\partial}{\partial x_j} \left(\overline{\theta'^2 u'_j} - \nu_\theta \frac{\partial \overline{\theta'^2}}{\partial x_j} \right) - 2\overline{\theta' u'_j \frac{\partial \bar{\theta}}{\partial x_j}} - 2\nu_\theta \overline{\frac{\partial \theta'}{\partial x_j} \frac{\partial \theta'}{\partial x_j}}. \quad (5.43)$$

乱れの運動量フラックスを支配する式

(5.33) $\times u'_k$ と、その i と k を置換した式を加え、全体の平均をとることにより求められる。(5.41) はこの式で $k = i$ の縮約をとったものである。

$$\begin{aligned}
\frac{\partial \overline{u'_i u'_k}}{\partial t} + \overline{u_j \frac{\partial \overline{u'_i u'_k}}{\partial x_j}} &= -\overline{u'_i u'_j \frac{\partial \bar{u}_k}{\partial x_j}} - \overline{u'_k u'_j \frac{\partial \bar{u}_i}{\partial x_j}} - \frac{\partial \overline{u'_i u'_k u'_j}}{\partial x_j} \\
- \frac{1}{\rho} \left[\frac{\partial \overline{p' u'_k}}{\partial x_i} + \frac{\partial \overline{p' u'_i}}{\partial x_k} - \overline{p' \left(\frac{\partial u'_i}{\partial x_k} + \frac{\partial u'_k}{\partial x_i} \right)} \right] &+ \nu \frac{\partial^2 \overline{u'_i u'_k}}{\partial x_j^2} - 2\nu \overline{\frac{\partial u'_i}{\partial x_j} \frac{\partial u'_k}{\partial x_j}}. \quad (5.44)
\end{aligned}$$