

$$\begin{aligned}
&= \frac{1}{(2\pi)^3} \int dr \mathbf{u} \cdot \nabla \left[\int d\mathbf{k}' \tilde{\mathbf{u}}(\mathbf{k}') \exp(i\mathbf{k}' \cdot \mathbf{r}) \right] \exp(-i\mathbf{k} \cdot \mathbf{r}) \\
&= \frac{1}{(2\pi)^3} \int dr \left[\int d\mathbf{k}' (\mathbf{u} \cdot i\mathbf{k}') \tilde{\mathbf{u}}(\mathbf{k}') \exp(i\mathbf{k}' \cdot \mathbf{r}) \right] \exp(-i\mathbf{k} \cdot \mathbf{r}) \\
&= \int d\mathbf{k}' \tilde{\mathbf{u}}(\mathbf{k}') \underbrace{i\mathbf{k}' \cdot \left[\frac{1}{(2\pi)^3} \int dr \mathbf{u} \exp[i(\mathbf{k}' - \mathbf{k}) \cdot \mathbf{r}] \right]}_{\tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')} \\
&= i \int d\mathbf{k}' \tilde{\mathbf{u}}(\mathbf{k}') (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}'))
\end{aligned}$$

となるので、ナビエ ストークス方程式の両辺をフーリエ変換したものは

$$\frac{\partial \tilde{\mathbf{u}}(\mathbf{k})}{\partial t} + i \int d\mathbf{k}' \tilde{\mathbf{u}}(\mathbf{k}') (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')) = -\frac{i}{\rho} \mathbf{k} \tilde{p} - \nu k^2 \tilde{\mathbf{u}}(\mathbf{k}). \quad (5.69)$$

(5.67) を用いて、この式からさらに \tilde{p} を消去することができる。(5.69) と \mathbf{k} の内積をとり、(5.67) を用いて $\mathbf{k} \cdot \tilde{\mathbf{u}}$ を 0 におくと

$$\frac{\partial}{\partial t} \underbrace{\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k})}_0 + i \int d\mathbf{k}' (\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k}')) (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')) = -\frac{i}{\rho} k^2 \tilde{p} - \nu k^2 \underbrace{\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k})}_0$$

より

$$-\frac{i}{\rho} \tilde{p} = \frac{i}{k^2} \int d\mathbf{k}' (\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k}')) (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')).$$

この \tilde{p} を (5.69) に代入して

$$\begin{aligned}
&\frac{\partial \tilde{\mathbf{u}}(\mathbf{k})}{\partial t} + i \int d\mathbf{k}' \tilde{\mathbf{u}}(\mathbf{k}') (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')) \\
&= i \int d\mathbf{k}' (\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k}')) (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')) \frac{\mathbf{k}}{k^2} - \nu k^2 \tilde{\mathbf{u}}(\mathbf{k}),
\end{aligned}$$

つまり

$$\left(\frac{\partial}{\partial t} + \nu k^2 \right) \tilde{\mathbf{u}}(\mathbf{k}) = -i \int d\mathbf{k}' (\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')) \left[\tilde{\mathbf{u}}(\mathbf{k}') - \frac{\mathbf{k}}{k^2} (\mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k}')) \right].$$

さらに (5.67) で波数を $\mathbf{k} - \mathbf{k}'$ とおくと

$$\mathbf{k}' \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}') = \mathbf{k} \cdot \tilde{\mathbf{u}}(\mathbf{k} - \mathbf{k}')$$